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Mai Ibrahim Khaleel,^{a,b,c} Yu-Da Chen,^{a,b,c} Ching-Hang Chien,^{a,b,c} and Yia-Chung Chang^{a,*}

^aAcademia Sinica, Research Center for Applied Sciences, Taipei, Taiwan ^bAcademia Sinica and National Tsing Hua University, Nano Science and Technology Program, Taiwan International Graduate Program, Taiwan ^cNational Tsing Hua University, Department of Engineering and System Science, Hsinchu, Taiwan

Abstract. Microscopic imaging ellipsometry is an optical technique that uses an objective and sensing procedure to measure the ellipsometric parameters Ψ and Δ in the form of microscopic maps. This technique is well known for being noninvasive and label-free. Therefore, it can be used to detect and characterize biological species without any impact. Microscopic imaging ellipsometry was used to measure the optical response of dried Streptococcus mutans cells on a glass substrate. The ellipsometric Ψ and Δ maps were obtained with the Optrel Multiskop system for specular reflection in the visible range ($\lambda = 450$ to 750 nm). The Ψ and Δ images at 500, 600, and 700 nm were analyzed using three different theoretical models with single-bounce, two-bounce, and multibounce light paths to obtain the optical constants and height distribution. The obtained images of the optical constants show different aspects when comparing the single-bounce analysis with the two-bounce or multibounce analysis in detecting S. mutans samples. Furthermore, the height distributions estimated by two-bounce and multibounce analyses of S. mutans samples were in agreement with the thickness values measured by AFM, which implies that the two-bounce and multibounce analyses can provide information complementary to that obtained by a single-bounce light path. @ 2017 Society of Photo-Optical Instrumentation Engineers (SPIE) [DOI: 10.1117/1 .JBO.22.5.056005]

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Introduction 1

Spectroscopic ellipsometry (SE) is a popular metrology technique used to determine thin film thickness and its optical constants. The technique allows us to reconstruct a nanoscale resolution image of the material surface by fitting the measured ellipsometric parameters psi (Ψ) and delta (Δ) with a suitable theoretical simulation.¹ Here Ψ is a measure of the absolute value of the ratio of the complex Fresnel reflection coefficients (or reflectivities), r_p and r_s , for p- and s-polarized lights, and Δ is the phase difference between r_p and r_s . Combing SE with microscopy leads to a powerful noncontact imaging technique known as microscopic imaging ellipsometry (MIE).² MIE enables a large area imaging with a spatial resolution of 1 μ m and subnanometer-scale z-resolution.^{2,3} In addition to the subnanoscale thickness resolution, it does not require a vacuum ambient nor the labeling of molecules, as compared to other detection techniques that require labeling-usually fluorescent or enzymatic such as fluorescence imaging.4-6 The labeling itself in such techniques can be highly heterogeneous, leading to nonquantitative results.⁶ MIE enables faster real-time measurements (in seconds) in comparison with other slow scanning techniques such as AFM that requires a few minutes to scan one area. This is especially useful for biological processes that occur within seconds rather than minutes.^{7,8} These unique features show the versatility of this imaging technique which had increased its usage in sensing and analyzing biological systems.

One of the pioneering applications of imaging ellipsometry was to detect antigen-antibody interactions using a single wavelength light source and fitting the measured data before and after interaction happens.9 Other applications of imaging ellipsometry were reported, including studies of biological reaction in cancer cells,¹⁰ detection of insulin-antibody binding on solid substrate,¹¹ detection of latent fingermarks,¹² as well as mapping of multilamellar films, hydrated membranes, and fluid domains.¹²

Streptococcus mutans is an aerobic, gram-positive bacterium with coccus shape found in the human oral cavity.^{14,15} It has been reported as the main etiological agent of dental caries and normal inhabitant of dental plaque.¹⁶ S. mutans is known for its ability to synthesize extracellular polymeric substance (EPS), which works as a matrix holding microbial cells together to form a biofilm that attaches to teeth surface. This matrix serves as the colony of bacteria and contributes to structure maintenance and antibiotic resistance of the biofilm.^{17,18} In the metabolism process, S. mutans produces aggressive substances such as acids, which cause a low pH environment in the mouth and lead to corrosion of enamel and the formation of cavities.18,19

In recent years, several research groups have developed methods for detecting and studying S. mutans in either laboratory settings or field trials. To our knowledge, there have been no attempts to detect S. mutans cells in the visible region using MIE. In this work, we performed MIE measurements of S. mutans cells grown on a glass substrate and dried in air. We assumed a structure consisting of air/bacterial cells/glass substrate to study the dielectric properties and height distribution.

^{*}Address all correspondence to: Yia-Chung Chang, E-mail: viachang@gate. sinica.edu.tw

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Fig. 1 Schematic diagram for rotating compensator ellipsometry.

The measured Ψ and Δ maps were fitted using three different optical models to detect bacterial cells and the surrounding dried culturing medium. In the first model, we considered only a single-bounce light path with the reflection from the interface between air and bacterial cells. In the second model, we considered the coherent superposition of the reflection from the interface between glass substrate and sample and the first reflection and called this the "two-bounce" model. In the third model, we included all reflections in the structure and called it the "multibounce" model. AFM Veeco Innova²⁰ tapping mode was used in this study to map the surface topography of S. mutans cells and culturing medium. The observed height distributions of cells and culturing medium on glass substrate by AFM were compared to the deduced values of height distribution by two-bounce and multibounce analyses. The glass substrate was chosen to grow S. mutans for two main reasons. The first one is that a series of studies have been done by using a glass substrate to study and reveal many details about S. mutans adherence and other oral bacteria.²¹ The second reason is that glass substrate is easy to calibrate, which makes the analyses simpler.

2 Materials and Methods

2.1 Cell Cultures

An *S. mutans* strain (ATCC 25175) was purchased from Creative Life Sciences (CMP), Taipei, Taiwan.²² *S. mutans* was incubated in brain-heart infusion broth (BHI) purchased from CMP, Taipei, Taiwan,²² shaken at 37°C overnight. Another BHI solution was prepared to grow bacterial cells on a glass substrate. The solution consists of 7.4-g BHI, 20-g sucrose, and 200-ml Millipore water.²³ Sucrose was added in order to enhance the rate of cell growth on a glass substrate.

2.2 Cell Sample Preparation

Glass slides (BRAND GMBH + CO KG, Wertheim, Germany)²⁴ were cut into 1×1 cm² pieces. They were cleaned through sonication in acetone for 30 min, then for 20 min in isopropanol solution at room temperature. After that glass pieces were washed and sonicated in Millipore water for 10 min and finally dried in an oven at 60°C. *S. mutans* was cultivated on glass substrates as described by Hu et al.²³ The cleaned glass substrates were immersed into a Petri dish containing 10 ml of freshly prepared BHI medium with sucrose. Then 2.5 ml of incubated *S. mutans* cells in BHI medium overnight was added to the

Petri dish contacting the immersed substrates, sealed by parafilm and incubated for 72 h at 37°C. The glass samples were then washed three times with Millipore water and placed in a new Petri dish to dehydrate naturally for 1 day, then measured by MIE and AFM. Bare glass substrates kept in culturing medium (BHI and sucrose) without *S. mutans* inoculated were used in reference experiments.

2.3 Multiskop Imaging Ellipsometry

The ellipsometric parameters Ψ and Δ of *S. mutans* cells on glass substrate were measured using an Optrel Multiskop system,²⁵ as shown in Fig. 1 which is modified to a rotating compensator ellipsometry. Converting the system from null measurements (which requires rotating both the polarizer and analyzer) to RCE allowed us to avoid the image-shift problem caused by the rotating analyzer (a Glan–Taylor prism²⁶) and to a faster extraction of the ellipsometric image.¹

The setup of our MIE consists of a laser arm, sample stage, and detector arm. The light beam comes from a Fianium WhiteLase short wavelength supercontinuum laser that covers wavelengths from 400 to 1000 nm with power >300 mW,²⁷ then passes through a polarizer fixed at 45 deg and a rotating compensator. To enhance the intensity of the incident light, a 10× lens was inserted before the light hits the spot of interest on the sample. The reflected or scattered light from the sample then goes through an $80 \times$ objective lens²⁸ that makes our system capable of probing an area as small as 1 μ m × 1 μ m. The light beam then passes through an analyzer mounted at 90 deg with respect to the polarizer, and finally is collected by a chargecoupled device (CCD) camera (Pixelfly qe),²⁹ with a pixel size equal to 164 nm. Measurements of ellipsometric parameters in this study were done using specular mode. In this mode, both the laser arm and detector arm are rotated in opposite directions with the use of a goniometer. The data collected by the CCD are the reflected light beam from the sample which is analyzed to determine the ellipsometric parameters Ψ and Δ . In RCE, the reflected light intensity as a function of compensator angle $\phi_{\rm c}$ is given as

$$I(\phi_{\rm c}) = A_0 + A_2 \cos(2\phi_{\rm c}) + B_2 \sin(2\phi_{\rm c}) + A_4 \cos(4\phi_{\rm c}) + B_4 \sin(4\phi_{\rm c}).$$
(1)

We take measurements of $I(\phi_c)$ for 36 values of ϕ_c starting from 5 deg and ending at 355 deg with 10 deg spacing. These



Fig. 2 Schematic diagram for light scattering paths from the sample under consideration: (a) singlebounce detection, (b) two-bounce detection, and (c) multibounce detection.

data allow us to determine the Fourier coefficients in Eq. (1), which can be used to evaluate Ψ and Δ .

2.4 Theory and Numerical Calculations

The phase difference between r_p and r_s (where r_p and r_s are reflectivity for p- and s-polarized light beams) and the ratio of their amplitudes can be quantified by the ellipsometric measurements as Δ and tan Ψ , respectively. The analysis of RCE signals was done to determine Ψ and Δ in terms of Fourier coefficients of the output light intensity as a function of the compensator angle.¹⁰

The fundamental equation of ellipsometry is then written as

$$\rho = \tan \Psi e^{i\Delta} = \frac{r_p}{r_s},\tag{2}$$

where r_p and r_s are complex reflectivities for p and s polarizations. For the sample under consideration, the reflectivities can be calculated from Fig. 2 using three different models.

The light detected by an image sensor is composed of light paths after a single bounce, two bounces, or multiple bounces.³⁰ For the single-bounce case, the coherent superposition of directly reflected light after a single reflection event with the sample is detected, while in the two- and multibounce cases the light is allowed to hit the interface between the sample and the substrate before reaching the sensor.³¹

In traditional methods of shape and reflection estimations such as time of flight and photonic mixer devices for opaque objects, the single-bounce measurement is commonly used, while the multibounce effect is ignored since it is considered as a source of noise.³¹ Recent studies of shape and reflection estimations have shown that two-bounce and multibounce light paths can provide information not achievable by using only single-bounce analysis. The two-bounce light paths were used for recovery of shape and reflection estimations,³¹ while multibounce light paths have been used to characterize flat and transparent thin films using a spectroscopic rotating polarizer-analyzer ellipsometer.³² To fully characterize the sample under consideration in this paper, we incorporated analyses of single-bounce, two-bounce, and multibounce models for the data obtained using RCE. The derivation of formulas relevant to the three models to be used is given in Appendix A.

2.4.1 Single-bounce light path

The single-bounce light path was used to detect bumpy areas, including *S. mutans* cells and the culturing medium on glass substrate. Since each grid (with 5 pixel average) of bumpy

area has a curved surface and size of 820 nm which is comparable to the wavelength, it is assumed that light going through the two-bounce and multibounce paths can be neglected [see Fig. 2(a)]. Therefore, light scattered from bumpy areas may be better described by the coherent sum of single-bounce light rays. By calculating r_p and r_s , we can determine the refractive index n_1 and the extinction coefficient k_1 of bumpy areas according to the following equation:

$$\varepsilon_1 = (n_1 + ik_1)^2 = \frac{(1-\rho)^2}{(1+\rho)^2} \sin \theta_0^2 \tan \theta_0^2 + \sin \theta_0^2, \quad (3)$$

where ε_1 is the complex refractive index of bumpy area.

2.4.2 Two-bounce light path

The model with a two-bounce light path is used to detect *S. mutans* cells and the dried culturing medium that represents a mixture of rough and nearly flat areas on the glass substrate. For the two-bounce light path shown in Fig. 2(b), the reflectivity is given as

$$r_j = r_{01j} + (1 - r_{01j}^2)r_{12j} \ e^{2i\beta},\tag{4}$$

where *j* stands for *s*-polarization and *p*-polarization, and β represents the phase factor for light going through the film given as

$$\beta = \frac{2\pi h n_1 \cos \theta_1}{\lambda},\tag{5}$$

where *h* and n_1 are the thickness and refractive index of the sample, respectively. Assuming the sample does not absorb light in the visible region (i.e., $k_1 = 0$), then we have

$$\ln |\gamma| = \ln |e^{2i\beta}| = \ln \left| \frac{r_{01p} - \rho r_{01s}}{\rho (1 - r_{01s}^2) r_{12s} - (1 - r_{01p}^2) r_{12p}} \right|$$

= 0. (6)

An in-house MATLAB[®] program³³ was used to solve Eq. (6) and find the refractive index and thickness of the cells and medium.

2.4.3 Multibounce light paths

The model with multibounce light paths is also used to analyze the data. Based on the results of Figs. 4(a) and 4(b), we expect this model will give similar results as a two-bounce model (for our case with AOI = 52 deg and height = 80 and 350 nm),



Fig. 3 *S. mutans* on a glass substrate after 72-h culture: (a) $20 \ \mu m \times 20 \ \mu m$ AFM topography image, (b) height profile of culturing medium (red, labeled 1) and *S. mutans* cells (green, labeled 2), and (c) height distribution of culturing medium and *S. mutans* cells.

since the higher order terms in the multibounce model appear to be small. The reflectivity of multibounce model is given as

$$r_j = \frac{r_{01j} + r_{12j} \ e^{2i\beta}}{1 + r_{01j}r_{12j} \ e^{2i\beta}},\tag{7}$$

where r_{01j} and r_{12j} for *s*-polarization and *p*-polarization have been derived in Ref. 32. Assuming that $k_1 = 0$ in the visible range, the final solution is written as

$$\ln|\gamma| = 0,\tag{8}$$

where γ is given in Eq. (20) of Appendix A. Our in-house MATLAB[®] program was also used to solve Eq. (8) and find the refractive index n_1 of the sample. Given n_1 , the thickness of the sample can be obtained from Eq. (5) in terms of β , n_1 , and $\cos \theta_1$ as

$$h = \frac{\beta \lambda}{2\pi n_1 \cos \theta_1}.$$
(9)

3 Results and Discussion

For MIE measurements, three samples of *S. mutans* were investigated, while only the results for two samples are reported here, since the third sample gives qualitatively the same results. A carbon tape was attached to the backside of the glass substrate of these samples in order to eliminate backside reflections of the glass substrate in our measurements. For AFM measurements, two samples were investigated using tapping mode to map the surface morphology of cells and culturing medium. All the samples investigated by AFM and MIE were prepared under the same conditions.

3.1 AFM Results

AFM tapping mode was used to morphologically characterize *S. mutans* cells cultured on a glass substrate for 72 h. All images were obtained with a commercial AFM (Veeco Innova) of ambient conditions. The OMCLO-AC160TS, R3 tip was used for tapping mode images.³⁴ Figure 3 shows AFM images of *S. mutans* cells and heights distribution on a glass substrate with spatial resolution of 78 nm/pixel.

The image shown in Fig. 3(a) shows the topography of *S. mutans* cells with different heights, while Fig. 3(b) shows the height profiles of *S. mutans* cells in red and culturing medium in green. Figure 3(c) shows the height distribution of *S. mutans* and culturing medium on glass substrate. The figure revealed that *S. mutans* cells are of height from 250 to 500 nm with maximum count at 350 nm, and a few stacked-up cells have heights of 600 to 650 nm. These heights correspond to the bright spots in Fig. 3(a). The height distribution of dried culturing medium (BHI and sucrose) on glass substrates in Fig. 3(c) shows height ranging from 10 to 200 nm with maximum count at 80 nm.

3.2 MIE-RCE Results

Before performing ellipsometry measurements of S. mutans cells on glass substrate, the Ψ and Δ spectra were measured for a bare glass substrate with a carbon tape attached to the backside in order to block the backside reflections of the glass substrate. Cauchy model³⁵ was used to fit Ψ and Δ spectra in the visible range (450 to 750 nm) and obtain the refractive index of glass substrate (1.549, 1.538, 1.532 at 500, 600, and 700 nm, respectively). MIE measurements of S. mutans samples cultured on a glass substrate for 72 h were done using specular mode at an angle of incidence (AOI) of 52 deg in the visible range (450 to 750 nm) and the analyses were done for three specific wavelengths. In order to overcome the signal noise in ellipsometry measurements, the values of Ψ and Δ of 5 \times 5 pixels are averaged to produce a mean value which is registered for the center pixel. We shift the center pixel by 2 pixels (in both xand y directions) and then take the average of the 5×5 pixels around it, and then move on. This leads to a smoothed image with a total area of 58 grids \times 58 grids, with each grid being 328 nm \times 328 nm in size (which contains signals average over the surrounding $820 \text{ nm} \times 820 \text{ nm}$ area). The averaging procedure was done by using an in-house MATLAB® program to produce smoothed images of the sample.

To show the difference among single-bounce, two-bounce, and multibounce analyses, $\tan \Psi$ and $\cos \Delta$ obtained for the three models are plotted as functions of n_1 in the interval (1, 3) as shown in Fig. 4. Here, $\tan \Psi = |r_p/r_s|$ and $\cos \Delta = \operatorname{Re}(r_p/r_s)/\tan \Psi$. In Fig. 4, $\tan \Psi$ and $\cos \Delta$ are plotted for three different wavelengths 500, 600, and 700 nm, except



Fig. 4 tan Ψ and cos Δ as functions of n_1 at AOI of 52 deg for (a) two-bounce model (symbols) and multibounce model (continuous curves) at height = 80 nm, and (b) two-bounce model (symbols) and multibounce model (continuous curves) at height = 350 nm. Single-bounce results are also included in (a) to compare with two- and multibounce models.

for the single-bounce case in Fig. 4(a) since the result is independent of wavelength. The height of the sample considered in Figs. 4(a) and 4(b) was set to 80 and 350 nm, where 80 nm is the average height of dried culturing medium and 350 nm is the average height of *S. mutans* revealed by AFM measurements in Fig. 3(c).

Figures 4(a) and 4(b) show that results obtained by twobounce and multibounce analyses overlap well for two different heights. This means that results of samples analyzed by twobounce and multibounce models will be very close. As a calibration, in Fig. 5 we show the results of Ψ and Δ images for culturing medium sample (without *S. mutans*) at a wavelength



Fig. 5 CCD, Ψ , Δ , n_1 , k_1 , and h images of BHI sample using single-bounce, two-bounce, and multibounce models at AOI = 52 deg. The frame size is 19 μ m × 19 μ m. The units of Ψ and Δ are degrees.

equal to 600 nm analyzed using single-bounce, two-bounce, and multibounce models.

We noticed that the topology of the n_1 image is very similar to the Ψ image, while the k_1 image is very similar to the Δ image. Comparing n_1 images with the corresponding Ψ images in Fig. 5, we conclude that the cyan color areas in n_1 images with n_1 values ranging from 1.25 (1.39) to 1.3 (1.41) in single-bounce (two-bounce and multibounce) analyses coincide with the areas in dark blue of Ψ images. The yellow areas of n_1 images match the light blue areas in Ψ images, while the dark blue spots in n_1 images match the red spots in Ψ images. In single-bounce analysis, the yellow-color areas in the n_1 image correspond to the blue areas in the k_1 image which has negative (unphysical) values. This could be caused by depolarization of scattered light or it is an indication that those areas are flat and not suitable for single-bounce analysis. Even with unphysical k_1 values in some areas, the n_1 values obtained by the singlebounce model are still fairly close to the results obtained by two-bounce and multibounce methods. The blue spots in the n_1 image correspond to the yellow and red spots in the k_1 image which has positive values, indicating possible absorption there.

The n_1 images obtained from two-bounce and multibounce analyses are very close (with difference less than 0.02 in all pixels). In the multibounce analysis, there are two roots for n_1 corresponding to the "+" and "-" signs in Eq. (20). It is interesting to note that the physical results for n_1 of the culturing medium are found only by adopting the "-" sign in Eq. (20) for this sample.

Comparing the k_1 image in single-bounce analysis with h images in two-bounce and multibounce analyses, we found that the blue stripes in the k_1 image match the red stripes in h images with height between 200 and 280 nm. Other areas correspond to k_1 near zero and height less than 40 nm, which are non-absobing based on single-bounce analysis. It is interesting to note that the height of the culturing medium revealed by the two-bounce and

multibounce analyses is either below 40 nm (flat areas) or above 200 nm (bumpy areas) and the topography of the *h* image is very close to the k_1 image and Δ image.

3.2.1 Single-bounce analysis

The single-bounce model was first used to analyze *S. mutans* with culturing medium on glass substrate. The left-side images of Figs. 6 and 7 show the CCD images of two measured areas (area 1 and area 2) of interest for *S. mutans* samples. The measured Ψ and Δ images at wavelengths of 500, 600, and 700 nm for the two different areas are shown in the second and third columns of Figs. 6 and 7.

 Ψ and Δ images at different wavelengths were used to calculate ρ in Eq. (2). The obtained value of ρ is substituted in Eq. (3), and our in-house MATLAB[®] program was used to solve this equation and find values of refractive index n_1 and extinction coefficient k_1 for the three wavelengths and show their images in the fourth and fifth columns of Figs. 6 and 7.

We noticed again that the topology of the n_1 image is very similar to the Ψ image, while the k_1 image is very similar to the Δ image. Comparing n_1 images with the corresponding Ψ images in Figs. 6 and 7 for the three wavelengths, we conclude that areas in blue color of n_1 images with n_1 ranging from 1.2 to 1.4 with a few (dark blue) spots less than 1.2 coincide with the blue areas in Ψ images.

The cyan area and green stripes (green and yellow stripes) in Fig. 6 (Fig. 7) for k_1 images have positive values between 0 and 0.2 which are physically reasonable, and they correspond to blue areas in Δ images with values between -180 deg and -160 deg. The corresponding n_1 values for these areas are between 1.6 and 2 which correspond to bumpy areas in culturing medium. There are a few spots in k_1 images with values larger than 0.2 which indicates strong light trapping of the bumpy structure. The areas in k_1 images with negative values (blue



Fig. 6 CCD, Ψ , Δ , refractive index n_1 , and extinction coefficient k_1 images of *S. mutans* sample for area 1 using single-bounce analysis at AOI = 52 deg. The frame size is 19 μ m × 19 μ m. The units of Ψ and Δ are degrees.



Fig. 7 CCD, Ψ , Δ , refractive index n_1 , and extinction coefficient k_1 images of *S. mutans* sample for area 2 using single-bounce analysis at AOI = 52 deg. The frame size is 19 μ m × 19 μ m. The units of Ψ and Δ are degrees.

to cyan) are attributed to strong depolarization effect and the corresponding values of n_1 there are between 1.2 and 1.6.

It should be noted that the depolarization effect due to scattering from *S. mutans* cells and rough areas of culturing medium has been neglected in this model. Thus, the n_1 and k_1 values obtained can be rather crude. The k_1 value obtained can be misleading, since it is very sensitive to the phase and the depolarization effect. On the other hand, this one-bounce analysis is quite effective in locating which area is bumpy and the distribution of *S. mutans* cells can be related to the n_1 distribution obtained.

3.2.2 Two-bounce analysis

For the two-bounce analysis, the calculated ρ from measured Ψ and Δ images were substituted in Eq. (6). This equation was solved for different wavelengths to find the refractive index. The third column of Figs. 8 and 9 show the n_1 images for the two samples obtained by solving Eq. (6) for wavelengths of 500, 600, and 700 nm.

Areas in cyan for n_1 images in Figs. 8 and 9 show n_1 values ranging from 1.2 to 1.4, and areas with light green and some



Fig. 8 CCD, Ψ , Δ , refractive index n_1 , and height distribution h images of S. mutans sample for area 1 using two-bounce analysis at AOI = 52 deg. The frame size is 19 μ m × 19 μ m. The units of Ψ and Δ are degrees.



Fig. 9 CCD, Ψ , Δ , refractive index n_1 , and height distribution h images of S. mutans sample for area 2 using two-bounce analysis at AOI = 52 deg. The frame size is 19 μ m × 19 μ m. The units of Ψ and Δ are degrees.

grids with yellow, orange, and red colors show values ranging from 1.4 to 2.6. Comparing results of n_1 images obtained from two-bounce and single-bounce analyses, we found that areas in cyan in Figs. 8 and 9 correspond to the same areas in blue color in Figs. 6 and 7, and n_1 values obtained by both models are similar, although the k_1 values there are negative (unphysical). This indicates that bacteria cells surrounded by culturing medium represented by the blue areas in n_1 images in Figs. 6 and 7 can still be detected by single-bounce analysis, and the unphysical values of k_1 there are due to depolarization effect. This is confirmed by the two-bounce analysis.

The two-bounce model can be used to analyze and sense the *S. mutans* cells in areas where the height distribution of cells is uniform, similar to height profile at 350 nm as shown in AFM results (green line) of Fig. 3(b). In addition, the refractive index values obtained from the two-bounce analysis can be used to calculate the height distribution of *S. mutans* cells surrounded by culturing medium according to Eq. (9). Thus, the two-bounce analysis is more versatile for the current application. The fifth columns of Figs. 8 and 9 show the obtained height (*h*) distribution for area 1 and area 2. The blue and cyan areas with *h* ranging from 40 to 200 nm in Figs. 8 and 9 correspond to areas filled with culturing medium. Areas with *h* above 200 nm (green, yellow, orange, and red) correspond to areas filled with cells, in agreement with the height distribution of *S. mutans* cells and dried culturing medium obtained by AFM in Fig. 3(c).

3.2.3 Multibounce analysis

The first two columns of Fig. 10 show images of n_1 and h distributions *in* area 1 for $\lambda = 500$ nm obtained by solving Eq. (8) with both "–" and "+" signs in Eq. (20), respectively. In both calculations, the dark blue area represents the unphysical zone (since n_1 obtained there is less than 1). We found that the two solutions are mutually exclusive with the unphysical zone obtained with "–" sign being the physical zone obtained with "+" sign, and vice versa. Furthermore, the physical values of

 n_1 range from 1.3 to 1.5 (red zone) in the first column and from 1.5 to 3 (mixed-color zone) in the second column. The corresponding heights in the first column ranges from 200 to 300 nm, while in the second column it ranges from 20 to 180 nm. So, the two solutions of the multibounce analysis can be used to distinguish regions of low index (high *h* value) and high index (low *h* value), which correspond to regions filled with high and low concentrations of *S. mutans* cells, respectively. We can combine the physical solutions for n_1 obtained with both signs, and the outcome is presented in the third column. The result is very similar to that obtained by the two-bounce analysis as shown in the first row of Fig. 8.

In the above analysis, we have neglected the effects of depolarization and stray light, which can lead to errors in our estimate of n_1 and h. Thus, the n_1 and height distributions deduced can be different for different wavelengths as shown in Figs. 6 and 7 as well as in Figs. 8 and 9. The height distribution analysis for different wavelengths of S. mutans and dried culturing medium in our study should be wavelength independent and the variation of n_1 in the range of wavelengths considered should also be negligible. Thus, this variation in results for different wavelengths gives an estimate of the error caused by depolarization effect. The error estimate of n_1 values obtained by single bounce and two bounce (same as multibounce) analyses are reported in Table 1, which shows a standard deviation in the 1% range. For the h distribution, since h varies significantly from pixel to pixel in the samples considered and the image resolution depends on wavelength, it is not straight-forward to estimate the error in h distribution by comparing results for different wavelengths. However, we can see that in areas where the height distribution is uniform, the results obtained become rather insensitive to the wavelength, while in bumpy areas the variation can be significant. The overall ranges and the topologies of hin Figs. 8 and 9 for different wavelengths are also similar. However, changes of distributions for wavelength from 500 to 700 nm are clearly visible.



Fig. 10 Multibounce analysis for refractive index n_1 obtained by using "+" and "-" signs in Eq. (20) and height (*h*) distribution images of *S. mutans* sample for area 1 at $\lambda = 500$ nm at AOI = 52 deg. The frame size is 19 μ m × 19 μ m. The units of Ψ and Δ are degrees.

Table 1 Standard deviation and mean value of n_1 obtained by singlebounce and two-bounce models for area 1 and area 2 at AOI = 52 deg.

	Single	bounce	Two bounce				
	n	1	r	0 ₁			
	$\bar{\sigma}$	$\overline{n_1}$	$\bar{\sigma}$	$\overline{n_1}$			
Area1	0.144	1.576	0.169	1.557			
Area2	0.13	1.59	0.161	1.567			

Finally, we give some estimate of the statistical errors in our measurements. We determine the average standard deviation $(\bar{\sigma})$ by repeating the same measurements five times for MIE measurements of a bare glass substrate with carbon tape attached to the backside. The mean value of Ψ and Δ maps ($\bar{\Psi}, \bar{\Delta}$) and average standard deviation ($\bar{\sigma}$) of this study are shown in Table 2. These statistical values were calculated for single-pixel maps and maps obtained by averaging over five surrounding pixels, since results reported in this study were based on Ψ and Δ maps averaged over five surrounding pixels explained in Sec. 3.2. As shown in Table 2, the average standard deviation ($\bar{\sigma}$) for the 5-pixel average map is about one half of the single-pixel map, thus the resulting MIE image will be less noisy.

Table 2 Standard deviation and mean value of Ψ and Δ measured by MIE for glass substrate, area 1, and area 2 at AOI = 52 deg.

		Glass substrate				Area1			Area2				
		Ψ		Δ		Ψ		Δ		Ψ		Δ	
		$\bar{\sigma}$	$\bar{\Psi}$	$\bar{\sigma}$	$\overline{ \Delta }$	$\bar{\sigma}$	$\bar{\Psi}$	$\bar{\sigma}$	$\overline{ \Delta }$	$\bar{\sigma}$	Ψ	$\bar{\sigma}$	$ \Delta $
500 nm	Single pixel	0.155	14.62	0.653	178.2	0.33	9.825	2.314	167.2	0.3	9.169	2.07	175.9
	5-pixel average	0.111	14.71	0.524	178.9	0.131	10.32	0.993	178.1	0.113	8.102	0.726	173.1
600 nm	Single pixel	0.251	14.32	0.563	178.8	0.346	10.77	2.1	168.9	0.337	8.3	2.03	168.1
	5-pixel average	0.15	14.36	0.468	179.1	0.106	10.33	0.708	174.5	0.107	8.659	0.646	177
700 nm	Single pixel	0.269	14.77	0.64	178.8	0.476	9.951	2.533	177.6	0.442	7.481	2.377	178.3
	5-pixel average	0.127	14.45	0.526	179	0.157	10.34	0.765	173	0.134	10.42	0.78	174

We can also estimate the error introduced in our Ψ and Δ values determined by using a finite number of compensator angles (ϕ_c) in our measurements. If there is no statistical error or systematic error, one should be able to determine the five Fourier coefficients in Eq. (1) with just five values of ϕ_c . Given the measurements at 36 values of ϕ_c , we can select them all (with 10 deg spacing), or adopt only 18 values at 20 deg spacing, or 12 values at 30 deg spacing. We then have three ways to determine the Ψ and Δ values. The deviation of these determined values from the mean gives us an estimate of the measurement error. The estimated average standard deviation $\bar{\sigma}$ for areas 1 and 2 of our sample with *S. mutans* cells are also listed in Table 2.

For single-pixel Ψ maps, error ranges from 1.8% for glass substrate to 4.7% for area 1 and 5.9% for area 2. In singlepixel Δ maps, error ranges from 0.3% for glass substrate to 1.4% for area 1 and 1.3% for area 2, indicating that 5-pixel averaged maps are more accurate and reproducible by MIE.

4 Conclusions

In this paper, we performed MIE measurements in the visible range to detect dried S. mutans cells and surrounding culturing medium on glass substrate. We found that two-bounce and multibounce lead to very close results for refractive index values and height distributions. Incorporating single-bounce, two-bounce, and multibounce analyses can reveal more information about the biological sample under investigation. The two-bounce and multibounce models provide the height distribution, while the single-bounce model provides additional information on the extinction coefficient or the depolarization effect. Results obtained in this work show the ability of MIE to detect biological structures and extract their average refractive index and height distribution. One of the major motivations of using MIE in biosensing is to gain more knowledge about the dielectric properties of biological systems for better understanding and further development of biomaterials. Moreover, it constitutes important driving forces in developments directed toward sensors and diagnostic applications for use in clinic medicine.⁴ Results obtained in this work show that MIE is a valuable detecting technique, which can differentiate bacteria types by examining their morphologies (including shapes and heights) and dielectric properties. Since it is a label-free, noninvasive sensing technique, it does not require any sample preparation, and it can perform measurements fairly fast (within seconds). Because of the above features, MIE can be used to study the interaction mechanism between a substrate and adsorbed biological layer via real-time measurements and analysis. The amount of information obtained from the analysis of MIE measurements can be extended to include the effects due to surface roughness and molecular orientation.⁴ Finally, comparing this sensing tool to other label-free and noninvasive imaging techniques such as Raman spectroscopy and IR/thermal imaging, MIE provides complementary information of the samples compared to Raman spectroscopy³⁶ and IR/thermal imaging.³⁷ Since Raman spectroscopy measures the vibrational modes of the sample, it is more sensitive to the chemical nature of the sample, while MIE is capable of detecting the distribution of cells and measuring their thickness and dielectric properties. The IR/thermal imaging technique detects the thermal radiation from the sample, which cannot achieve the submicron spatial resolution considered here, and it does not have the capability of extracting detailed structure information as MIE does. Although we report MIE technique only in the visible range, the same experimental setup and analysis can be applicable for other wavelengths such as ultraviolet or infrared.

Appendix A: Derivation of Formulas for the Three Models Used

A.1 Single-Bounce Light Path

The Fresnel reflection coefficients r_p and r_s in z-direction for single-bounce light path from Fig. 2(a) are

$$r_{01p} = \frac{E_{rx}}{E_{ix}} = \frac{\frac{1}{\epsilon_0} K_{iz} - \frac{1}{\epsilon_1} K_{tz}}{\frac{1}{\epsilon_0} K_{iz} + \frac{1}{\epsilon_1} K_{tz}} = \frac{\frac{1}{\epsilon_0} \sqrt{\epsilon_0 - \sin \theta_0^2} - \frac{1}{\epsilon_1} \sqrt{\epsilon_1 - \sin \theta_0^2}}{\frac{1}{\epsilon_0} \sqrt{\epsilon_0 - \sin \theta_0^2} + \frac{1}{\epsilon_1} \sqrt{\epsilon_1 - \sin \theta_0^2}},$$
 (10)

$$r_{01s} = \frac{E_{ry}}{E_{iy}} = \frac{K_{0z} - K_{1z}}{K_{0z} + K_{1z}} = \frac{\sqrt{\varepsilon_0 - \sin\theta_0^2} - \sqrt{\varepsilon_1 - \sin\theta_0^2}}{\sqrt{\varepsilon_0 - \sin\theta_0^2} + \sqrt{\varepsilon_1 - \sin\theta_0^2}},$$
(11)

where K_{iz} (i = 0, 1) denotes the z-component of the wavevector for light in media i. θ_0 , ε_0 , and ε_1 are the AOI, the complex refractive index of air, and the complex refractive index of *S. mutans* cells, respectively. Substituting for r_p and r_s in Eq. (2) in Sec. 2.4, we obtain

$$\rho = \tan \Psi e^{i\Delta} = \frac{\left(\frac{1}{\varepsilon_0}\sqrt{\varepsilon_0 - \sin\theta_0^2} - \frac{1}{\varepsilon_1}\sqrt{\varepsilon_1 - \sin\theta_0^2}\right)\left(\sqrt{\varepsilon_0 - \sin\theta_0^2} + \sqrt{\varepsilon_1 - \sin\theta_0^2}\right)}{\left(\frac{1}{\varepsilon_0}\sqrt{\varepsilon_0 - \sin\theta_0^2} + \frac{1}{\varepsilon_1}\sqrt{\varepsilon_1 - \sin\theta_0^2}\right)\left(\sqrt{\varepsilon_0 - \sin\theta_0^2} - \sqrt{\varepsilon_1 - \sin\theta_0^2}\right)}.$$
(12)

Solving Eq. (12) with respect to ε_1 , we have

$$\varepsilon_1 = (n_1 + ik_1)^2 = \frac{(1-\rho)^2}{(1+\rho)^2} \sin \theta_0^2 \tan \theta_0^2 + \sin \theta_0^2.$$
 (13)

A.2 Two-Bounce Light Path

For the two-bounce light path shown in Fig. 2(b), the Fresnel reflection coefficients is given by

$$r_j = r_{01j} + (1 - r_{01j}^2) r_{12j} e^{2i\beta},$$
(14)

where *j* stands for *s*-polarization and *p*-polarization, and β represents the phase film factor given as

$$\beta = \frac{2\pi h n_1 \cos \theta_1}{\lambda},\tag{15}$$

where h and n_1 are the thickness and refractive index of the sample, respectively.

Substituting Eq. (14) for s and p polarizations in Eq. (2), Sec. 2.4, we obtain

$$\rho = \tan \Psi e^{i\Delta} = \frac{r_{01p} + (1 - r_{01p}^2)r_{12p}\gamma}{r_{01s} + (1 - r_{01s}^2)r_{12s}\gamma},$$
(16)

where r_{01p} , r_{01s} , r_{12p} , and r_{12s} have been derived in a previous work by Taya et al.³²

Solving Eq. (16) with respect to $e^{2i\beta}$ we obtain

$$\gamma = e^{2i\beta} = \frac{r_{01p} - \rho r_{01s}}{\rho (1 - r_{01s}^2) r_{12s} - (1 - r_{01p}^2) r_{12p}}.$$
(17)

Assuming the sample does not absorb light in the visible region (i.e., $k_1 = 0$), then

$$\ln|\gamma| = \ln \left| \frac{r_{01p} - \rho r_{01s}}{\rho (1 - r_{01s}^2) r_{12s} - (1 - r_{01p}^2) r_{12p}} \right| = 0.$$
(18)

A.3 Multibounce Light Path

The Fresnel reflection coefficients of multibounce light paths have been derived in Ref. 32 and given as

$$\rho = \tan \Psi e^{i\Delta} = \frac{(r_{01p} + r_{12p}\gamma)}{(r_{01s} + r_{12s}\gamma)} \frac{(1 + r_{01s}r_{12s}\gamma)}{(1 + r_{01p}r_{12p}\gamma)}.$$
 (19)

Assuming that $k_1 = 0$ in the visible range, then the final solution is written as

$$\ln|\gamma| = \ln|e^{2i\beta}| = \ln\left|\frac{-B \pm \sqrt{B^2 - 4AC}}{2A}\right| = 0,$$
 (20)

where

$$A = (\rho r_{01p} - r_{01s})(r_{12p}r_{12s}), \tag{21}$$

$$B = r_{01p}r_{01s}(\rho r_{12p} - r_{12s}) + \rho r_{12s} - r_{12p}, \qquad (22)$$

$$C = \rho r_{01s} - r_{01p}. \tag{23}$$

The thickness of the sample can be obtained from Eq. (15) in terms of β , n_1 , and $\cos \theta_1$ as

$$h = \frac{\beta \lambda}{2\pi n_1 \cos \theta_1}.$$
(24)

Disclosures

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Mai Ibrahim Khaleel is a PhD student in the nanoscience and technology program—Taiwan International Graduate Program—in collaboration with the Department of Engineering and System Science at National Tsing Hua University. She received her BS and MS degrees in physics from Al-Quds University in 2007 and 2009, respectively. Her current research interests include bio-optics, imaging ellipsometry, and photonics.

Biographies for the other authors are not available.